

# Fluktuation - Dissipationstheorem

Wickel - Formel  $V = - \int d\vec{r}^3 F(\vec{r}, t) \hat{A}(\vec{r})$

$$\langle \delta B(\vec{r}, t) \rangle = \int dt' d\vec{r}' \chi(\vec{r} - \vec{r}', t - t') F(\vec{r}', t')$$

$$\begin{aligned} \chi_{BA}(\vec{r} - \vec{r}', t - t') &= \frac{i}{\hbar} \text{Tr} \left\{ \rho_0 [B_I(\vec{r}, t), A_I(\vec{r}', t')] \right\} \Theta(t - t') \\ &= \frac{i}{\hbar} \langle [B(\vec{r}, t), A(\vec{r}', t')] \rangle \Theta(t - t') \end{aligned}$$

~~Resonanz~~  $\rho_0 = \frac{1}{Z} e^{-\frac{H_0}{kT}} = \frac{1}{Z} e^{-\beta H_0}$

Operatoren im Gleichgewicht

$$B = B_T = B_H \quad \text{wenn} \quad H = H_0$$

Fall  $A = B$   $H_0 |n\rangle = E_n |n\rangle$ ,  $|n\rangle$  vollständig

$$\chi_{AA} = \frac{i}{\hbar} \frac{\Theta(t - t')}{Z_0} \sum_n \langle n | e^{-B H_0} [A(t), A(t')] | n \rangle$$

$$A(t) = e^{i \frac{H_0 t}{\hbar}} A_S e^{-i \frac{H_0 t}{\hbar}}$$

$$\begin{aligned} &= \frac{i}{\hbar} \frac{\Theta(t - t')}{Z_0} \sum_m \left( \langle m | e^{-B H_0} e^{i \frac{H_0 t}{\hbar}} A_S(m) e^{-i \frac{H_0 t}{\hbar}} e^{i \frac{H_0 t'}{\hbar}} A_S(r') \right. \\ &\quad \left. e^{-i \frac{H_0 t'}{\hbar}} | m \rangle \right. \\ &\quad \left. - \langle r \leftrightarrow r' \rangle \right) \sum_{m'} | m' \rangle \langle m' | \end{aligned}$$

$$= \frac{i}{\hbar} \frac{\Theta(t - t')}{Z_0} \sum_{m, m'} \langle m | e^{-\beta E_m} e^{i \frac{E_m t}{\hbar}} A_S(r) e^{-i \frac{E_m t}{\hbar}} | m' \rangle$$

$$\langle m' | e^{i \frac{E_m t'}{\hbar}} A_S(r') e^{-i \frac{E_m t'}{\hbar}} | m \rangle - \langle r \leftrightarrow r' \rangle$$

$$= \frac{i}{\hbar} \frac{\Theta(t-t')}{Z_0} \sum_m \langle m | A(r) | m' \rangle \langle m' | A(r') | n \rangle e^{-\beta E_n \frac{\hbar}{k} T_n(t-t')} e^{-\frac{i}{\hbar} E_n (t-t')}$$

~  
 $r \leftrightarrow r'$   
 $t \leftrightarrow t'$

FT  $t \rightarrow \omega$

$$\int_0^\infty dt e^{i\omega t} e^{-\delta t} = \frac{1}{i\omega + \delta} = i \frac{1}{\omega + i\delta}$$

$$\underset{\delta \rightarrow 0}{=} i \left( PV \frac{1}{\omega} - i\pi \delta(\omega) \right)$$

$$\chi(r, r', \omega) = \int_{-\infty}^{\infty} d\tau \chi(r, r', \tau) e^{i\omega \tau}$$

$\tau = t - t'$

$$= \frac{i}{\hbar} \frac{1}{Z_0} \sum_{m, m'} \langle m | A(r) | m' \rangle \langle m' | A(r') | n \rangle \frac{i e^{-\beta E_n}}{\omega + \frac{E_m - E_{m'}}{\hbar} + i\delta}$$

$$- \langle n | A(r') | m' \rangle \langle m' | A(r) | n \rangle \frac{i e^{-\beta E_n}}{\omega + \frac{E_{m'} - E_m}{\hbar} + i\delta}$$

$$\chi(k, \omega) = \int d^3(r-r') \chi(r-r', \omega) e^{-ik(r-r')}$$

wahle  $r = r' \Rightarrow A(r) = A(r') = A$

$$\chi(\omega) = -\frac{1}{\hbar Z_0} \sum_{m, m'} |\langle m | A | m' \rangle|^2 e^{-\beta E_n}$$

$$\left( \frac{1}{\omega + \frac{E_m - E_{m'}}{\hbar} + i\delta} - \frac{1}{\omega - \frac{E_m - E_{m'}}{\hbar} + i\delta} \right)$$

Def.  $\chi = \chi' + i\chi'' \quad \chi' = \text{Re}(\chi) \quad \chi'' = \text{Im}(\chi)$

$$\chi'' = \frac{\pi}{\hbar Z_0} \sum_{m, m'} |\langle m | A | m' \rangle|^2 \left\{ \delta\left(\omega + \frac{E_m - E_{m'}}{\hbar}\right) - \delta\left(\omega - \frac{E_m - E_{m'}}{\hbar}\right) \right\} e^{-\beta E_n}$$

$$\begin{aligned}
&= \frac{\pi}{kZ_0} \left( \sum_{m, m'} |\langle n | A | m' \rangle|^2 \delta(\omega + \frac{E_m - E_{m'}}{k}) e^{-\beta E_m} \right. \\
&\quad \left. + \sum_{m', m} |\langle n' | A | m \rangle|^2 \delta(\omega + \frac{E_{m'} - E_m}{k}) e^{-\beta E_{m'}} \right) \\
&= \frac{\pi}{kZ_0} \sum_{m, m'} |\langle n | A | m' \rangle|^2 \delta(\omega + \frac{E_m - E_{m'}}{k}) (e^{-\beta E_m} - e^{-\beta E_{m'}})
\end{aligned}$$

$$\frac{E_m - E_{m'}}{k} = -\omega \Rightarrow E_{m'} = E_m + k\omega$$

$$X = \frac{\pi e^{-\beta E_m} (1 - e^{-\beta k\omega})}{k Z_0} \sum_{m, m'} |\langle m | A | m' \rangle|^2 \delta(\omega + \frac{E_m - E_{m'}}{k})$$

$$\langle A(t) A(t') \rangle_o = \text{Tr} (\rho A(t) A(t')) \quad \text{Korrelationsfunktion}$$

$$S(t-t') = \frac{1}{2} \langle A(t) A(t') + A(t') A(t) \rangle_o$$

Symmetrisierte Korrelationsfunktion - Rauschen

$$\langle A \rangle = 0 \quad \text{aber } \overbrace{A}^{\text{wie charakterisiert man } A?} \text{ (das Rauschen)}$$

wie charakterisiert man  $A \in \mathbb{C}$  (das Rauschen)

$\sigma = |t-t'|$  Korrelation klein, wenn  
 $A(t')$  unabhängig von  $A(t)$

$$\begin{aligned}
S(t-t') &= \frac{1}{2Z_0} \sum_n \left\langle n | e^{-\beta H_0} e^{i \frac{H_0 t}{k}} A e^{-i \frac{H_0 t'}{k}} e^{\frac{i}{k} H_0 t'} A e^{-\frac{i}{k} H_0 t} | n \right\rangle \\
&= \frac{1}{2Z_0} \sum_{m, m'} |\langle m | A | m' \rangle|^2 e^{-\beta E_m} e^{\frac{i}{k} (E_m - E_{m'}) (t - t')} + t \leftrightarrow t'
\end{aligned}$$

$$FT$$

$$S(\omega) = \frac{1}{2\pi} \sum_{n, m'} |\langle n | A | m' \rangle|^2 e^{-\beta E_m} \left[ 2\pi \left( \delta\left(\omega + \frac{E_m - E_{m'}}{\hbar}\right) + \delta\left(\omega - \frac{E_m + E_{m'}}{\hbar}\right) \right) \right]$$

$$\int dt e^{i\omega t} = 2\pi \delta(\omega)$$

$$= \frac{\pi}{Z_0} e^{-\beta E_n} (1 + e^{-\beta \hbar \omega}) \sum_{n, m'} |\langle n | A | m' \rangle|^2 \delta\left(\omega + \frac{E_m - E_{m'}}{\hbar}\right)$$

$$S(\omega) = \hbar \coth\left(\frac{\beta \hbar \omega}{2}\right) \chi''(\omega)$$

Klassische Lösung

$$\coth(x) \approx \frac{1}{x} \quad x \ll 1$$

$$\text{für } k_B T \gg \hbar \omega$$

$$S(\omega) \approx \frac{2}{\omega_B} \chi''(\omega) = \frac{2k_B T}{\omega} \chi''(\omega)$$

Quantensprung  $\longleftrightarrow$  Klass. Rauschen

$$\hbar \omega \gtrsim k_B T$$

$$\hbar \omega \ll k_B T$$

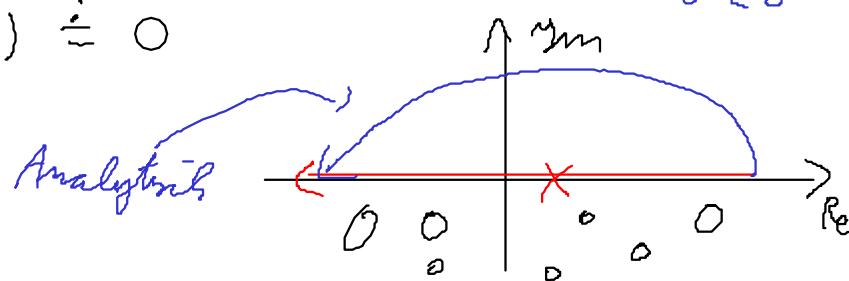
$\chi(t-t') \propto \Theta(t-t') \Rightarrow \chi(\omega)$  Analytisch

$$\text{für } \gamma_m(\omega) > 0$$

$$\chi(\tau) = \int \frac{d\omega}{2\pi} e^{-i\omega \tau} \chi(\omega) \quad \omega = i\alpha \Rightarrow e^{\alpha \tau}$$

$$\text{mit } \chi(\tau < 0) \stackrel{!}{=} 0$$

$$\tau < 0$$



$$\chi'(\omega) = \frac{PV}{\pi} \int d\omega' \frac{\chi''(\omega')}{\omega' - \omega}$$

Kramers  
Kronig  
Relation

$$\chi''(\omega) = - \frac{PV}{\pi} \int d\omega' \frac{\chi'(\omega')}{\omega' - \omega}$$